ГОДИШНИК НА СОФИЙСКИЯ УНИВЕРСИТЕТ "СВ. КЛИМЕНТ ОХРИДСКИ" ФАКУЛТЕТ ПО МАТЕМАТИКА И ИНФОРМАТИКА Том 96

ANNUAIRE DE L'UNIVERSITE DE SOFIA "ST. KLIMENT OHRIDSKI"
FACULTE DE MATHEMATIQUES ET INFORMATIQUE
Tome 96

ESTIMATES FOR THE SINGULAR SOLUTIONS OF THE 3-D PROTTER'S PROBLEM

NEDYU POPIVANOV and TODOR POPOV

For the wave equation we study boundary value problems, stated by Protter in 1952, as some three-dimensional analogues of Darboux problems on the plane. It is known that Protter's problems are not well posed and the solution may have singularity at the vertex O of a characteristic cone, which is a part of the domain's boundary $\partial\Omega$. It is shown that for n in $\mathbb N$ there exists a right-hand side smooth function from $C^n(\bar{\Omega})$, for which the corresponding unique generalized solution belongs to $C^n(\bar{\Omega}\setminus O)$, but it has a strong power-type singularity. It is isolated at the vertex O and does not propagate along the cone. The present article gives some necessary and sufficient conditions for the existence of a fixed order singularity. It states some exact a priori estimates for the solution.

Keywords: wave equation, boundary value problems, generalized solution, singular solutions, propagation of singularities, special functions

2000 MSC: main 35L05, 35L20, 35D05, 35A20; secondary 33C05, 33C90

1. INTRODUCTION

We discuss some boundary value problems for the wave equation

$$\Box u = u_{x_1 x_1} + u_{x_2 x_2} - u_{tt} = f \tag{1.1}$$

in a simply connected domain $\Omega \subset \mathbb{R}^3$. The domain

$$\Omega := \{ (x_1, x_2, t) : 0 < t < 1/2, \ t < \sqrt{x_1^2 + x_2^2} < 1 - t \}$$

is bounded by two characteristic cones of (1.1)

$$\begin{split} S_1 &= \{(x_1, x_2, t): \ 0 < t < 1/2, \ \sqrt{x_1^2 + x_2^2} = 1 - t\}, \\ S_2 &= \{(x_1, x_2, t): \ 0 < t < 1/2, \ \sqrt{x_1^2 + x_2^2} = t\} \end{split}$$

and the circle $S_0 = \{(x_1, x_2, t) : t = 0, x_1^2 + x_2^2 < 1\}$ centered at the origin O(0, 0, 0). The following three-dimensional analogues of the plane Darboux problems are stated by M. Protter [27]:

Problems P1 and P2. To find a solution of the wave equation (1.1) in Ω , which satisfies one of the following boundary conditions:

P1
$$u|_{S_0} = 0$$
 and $u|_{S_1} = 0$;
P2 $u_t|_{S_0} = 0$ and $u|_{S_1} = 0$.

The corresponding adjoint problems are:

Problems P1* and P2*. To find a solution of the wave equation (1.1) in Ω , which satisfies the corresponding boundary conditions:

P1*
$$u|_{S_0} = 0$$
 and $u|_{S_2} = 0$;
P2* $u_t|_{S_0} = 0$ and $u|_{S_2} = 0$.

For the recent known results concerning Protter's problems see [25] and references therein. For further publications in this area see [1, 2, 8, 13, 16, 19, 20].

Substituting the boundary condition on S_0 by $[u_t + \alpha u]|_{S_0} = 0$, one obtains Problem P_{α} , for which we refer to [11] and references therein. In the case of the wave equation, involving either lower order terms or some other type perturbations, Problem P2 in Ω has been studied in [1, 2, 3, 12]. On the other hand, Bazarbekov [5] gives another analogue of the classical Darboux problem in the same domain Ω . Some other statements of Darboux type problems can be found in [4, 6, 18] in bounded or unbounded domains different from Ω .

Protter [27] formulated and studied these three-dimensional analogues of the Darboux problem on the plane 50 years ago – in 1952. Nowadays, it is known that in contrast to the Darboux problem in \mathbb{R}^2 the 3 – D Problems P1 and P2 are not well posed. The reason for this is that the adjoint homogeneous Problems $P1^*$ and $P2^*$ have smooth solutions and the linear space they generate is infinite dimensional as one could see in Tong Kwang-Chang [29], Popivanov, Schneider [24], Khe Kan Cher [20] and Popivanov, Popov [26].

Lemma 1.1. (see [13]) Let ρ , φ and t be the polar coordinates in \mathbb{R}^3 : $x_1 = \rho \cos \varphi$, $x_2 = \rho \sin \varphi$. Let us define the functions

$$H_k^n(\rho,t) := \sum_{i=0}^k A_i^k \frac{t(\rho^2 - t^2)^{n-3/2 - k - i}}{\rho^{n-2i}}$$
 (1.2)

and

$$E_k^n(\rho,t) = \sum_{i=0}^k B_i^k \frac{(\rho^2 - t^2)^{n-1/2 - k - i}}{\rho^{n-2i}},$$

where the coefficients are

$$A_i^k := (-1)^i \frac{(k-i+1)_i (n-1/2-k-i)_i}{i!(n-i)_i}$$

and

$$B_i^k := (-1)^i \frac{(k-i+1)_i (n+1/2-k-i)_i}{i!(n-i)_i}$$

with $(a)_i := a(a+1)...(a+i-1)$. Then for $n \in \mathbb{N}$, $n \geq 4$, the functions

$$V_k^{n,1}(\rho,t,\varphi) = H_k^n(\rho,t)\cos n\varphi \qquad and \qquad V_k^{n,2}(\rho,t,\varphi) = H_k^n(\rho,t)\sin n\varphi$$

are classical solutions of the homogeneous Problem P1* for all $k = 0, 1, ..., \left[\frac{n}{2}\right] - 2$, and the functions

$$W_k^{n,1}(\rho,t,\varphi) = E_k^n(\rho,t)\cos n\varphi \qquad and \qquad W_k^{n,2}(\rho,t,\varphi) = E_k^n(\rho,t)\sin n\varphi$$

are classical solutions of the homogeneous Problem P2*

for all
$$k = 0, 1, ..., \left[\frac{n-1}{2}\right] - 1$$
.

A necessary condition for the existence of classical solution for Problem P1 (Problem P2) is the orthogonality of the right-hand side function f to all functions $V_k^{n,i}(\rho,t,\varphi)$ (respectively $W_k^{n,i}$). To avoid an infinite number of necessary conditions in the frame of classical solvability, we need to introduce some generalized solutions of Problems P1 and P2 with eventually singularity on the characteristic cone Σ_2 , or only at its vertex O. Popivanov, Schneider in [24] and [25] give the following definition:

Definition 1.1. A function $u = u(x_1, x_2, t)$ is called a generalized solution of the Problem P1 in Ω if:

- 1) $u \in C^1(\bar{\Omega} \setminus O), u|_{S_0 \setminus O} = 0, u|_{S_1} = 0;$
- 2) the identity

$$\int\limits_{\Omega} (u_t w_t - u_{x_1} w_{x_1} - u_{x_2} w_{x_2} - f w) \, dx_1 dx_2 dt = 0$$

holds for all $w \in C^1(\bar{\Omega})$, w = 0 on S_0 , and w = 0 in a neighborhood of S_2 .

Garabedian [10] proved the uniqueness of a classical solution of Problem P1. Popivanov, Schneider [25] proved the uniqueness of a generalized solution in $C^1(\overline{\Omega}\backslash O)$. It is known (cf. Popivanov, Schneider [25], Aldashev [1]) that for every $n \in \mathbb{N}$, $n \geq 4$, there exists a smooth function

$$f_n(x_1, x_2, t) := f_n^1(\rho, t) \cos n\varphi \in C^{n-2}(\overline{\Omega}),$$

whose corresponding generalized solution of Problem P1 near the origin O behaves like r^{1-n} , where $r = (x_1^2 + x_2^2 + t^2)^{1/2}$ is the distance to the origin. The same phenomenon appears in the case of a more general boundary condition P_{α} (see [11]). These singularities of the generalized solutions do not propagate in the direction of the bicharacteristics on the characteristic cone. It is traditionally assumed that

the wave equation, whose right-hand side is sufficiently smooth in $\bar{\Omega}$, cannot have a solution with an isolated singular point. For results concerning the propagation of singularities for operators of second order we refer to Hörmander [15, Chapter 24.5]. For some related results in the case of the plane Darboux problem see [23].

Further, we study the Problem P1 only in the case when the right-hand side function f is a trigonometric polynomial of order l:

$$f(x_1, x_2, t) = f_0^1(\rho, t) + \sum_{n=1}^{l} \left(f_n^1(\rho, t) \cos n\varphi + f_n^2(\rho, t) \sin n\varphi \right). \tag{1.3}$$

In this case Popivanov, Schneider [25] proved the existence and uniqueness of a generalized solution $u(x_1, x_2, t)$ of the corresponding problem. We already know that $u(x_1, x_2, t)$ may have power type singularity at the origin. More precisely, there are solutions that have the growth of r^{1-l} at the point O. In this paper we will prove some existence and uniqueness results for Problems P1 and P2 and study the behavior of the generalized solution around the origin. Let us denote the weighted uniform norm

$$||f_n||_q^* := \sum_{|\alpha| \le q: \ i=1,2} \max_{0 \le t \le \rho \le 1-t} \left| r^{|\alpha|+1/2} D^{\alpha} f_n^i(\rho, t) \right|,$$

analogous to the weighted Sobolev norms in corner domains (see [21], [14]). Denote as usual $x_+^{\alpha} = x^{\alpha}$ for x > 0, and $x_+^{\alpha} = 0$ for $x \leq 0$. Then the main results are:

Theorem 1.1. Let us suppose that $f(x_1, x_2, t) \in C^{(l-4)_+}(\overline{\Omega})$ has the form (1.3) and that

$$\int_{\Omega} V_k^{n,i}(x_1, x_2, t) f(x_1, x_2, t) dx_1 dx_2 dt = 0$$
(1.4)

for all n=2,3,...,l; i=1,2; $k=0,1,...,\left[\frac{n}{2}\right]-1.$ Then there exists an unique generalized solution u of Problem P1 . Moreover, $u\in C^{1+(l-4)+}(\overline{\Omega}\backslash O)$ and for every $\varepsilon>0$ it holds the a priori estimate

$$|u(x_{1}, x_{2}, t)| \leq C_{1} r^{1/4} ||f_{0}||_{0}^{*} + C_{2, \varepsilon} r^{-\varepsilon} ||f_{1}||_{0}^{*}$$

$$+ C_{3} r^{1/2} |\ln r| \sum_{k=1}^{\lfloor l/2 \rfloor} ||f_{2k}||_{(2k-4)_{+}}^{*}$$

$$+ C_{4} \sum_{k=1}^{\lfloor \frac{l-1}{2} \rfloor} ||f_{2k+1}||_{(2k-3)_{+}}^{*},$$

$$(1.5)$$

where $C_{2,\varepsilon}$ depends on ε , but all the constants C_1 , $C_{2,\varepsilon}$, C_3 and C_4 are independent on the function f.

Theorem 1.1 gives an a priori estimate of the generalized solution. Now the next Theorem 1.2 provides an a priori estimate for the generalized solution and clarifies the significance of the above orthogonality conditions (1.4). In other worlds, for any couple (n, k) the corresponding condition "controls" one power-type singularity.

Theorem 1.2. Let $s \in \mathbb{N}$ be such that $2 \le s \le l$. Suppose that $f(x_1, x_2, t) \in C^{(l-1)+}(\overline{\Omega})$ has the form (1.3) and satisfies the orthogonality conditions (1.4) for any couple (n,k) such that $n,k \in \mathbb{N} \cup \{0\}$, $2 \le n \le l$, $n-2k \ge s+1$ and i=1,2. Then there exists an unique generalized solution u of Problem P1 such that: $u \in C^{1+(l-4)+}(\overline{\Omega}\backslash O)$ and the estimate

$$|u(x_1, x_2, t)| \le Cr^{-(s-1)} \sum_{k=0}^{l} ||f_k||_{(k-s-1)_+}^*$$
 (1.6)

holds. If we suppose additionally that there are $m, p \in \mathbb{N} \cup \{0\}$ and j = 1 or 2 such that $2 \le m \le l$, m - 2p = s and

$$\int_{\Omega} V_p^{m,j}(x_1, x_2, t) f(x_1, x_2, t) dx_1 dx_2 dt \neq 0, \tag{1.7}$$

then in some neighborhood of the origin one has

$$\left|u_m^j\right|_{S_2}(x_1, x_2)\right| \ge c(x_1^2 + x_2^2)^{-(s-1)/2}, \quad c > 0,$$
 (1.8)

where

$$u_m^1|_{S_2}(x_1,x_2) := \int\limits_0^{2\pi} u|_{t=|x|} \cos m\varphi \; d\varphi, \quad u_m^2|_{S_2}(x_1,x_2) := \int\limits_0^{2\pi} u|_{t=|x|} \sin m\varphi \; d\varphi.$$

To illustrate the dependence of the singularity of the generalized solution on the orthogonality assumptions, let us consider the following table:

Table 1. The orthogonality conditions and the order of singularity

	l	l-1	1-2	l - 3	 m		***	4	3	2
1					 			$V_1^{4,i}$	♦ 173,i	$V_0^{2,i}$
2 3					 			$V_0^{4,i}$	$\overset{\diamond}{V_0^{3,i}}$	
m-2p-1					 $V_p^{m,i}$	***				
m - 2p - 1 $l - 4$ $l - 3$ $l - 2$ $l - 1$	$V_1^{l,i}$	$V_1^{l-1,i} \\ \diamond \\ V_2^{l-1,i}$	$V_0^{l-2,i}$	$V_0^{l-3,i}$	 •••					
l-1	$V_0^{l,i}$,0								

Observe that both $V_k^{n,1}$ and $V_k^{n,2}$ are located in the *n*-th column and (n-2k-1)-st row of Table 1. Thus, $V_0^{n,i}$ form the right most diagonal, the next one is empty – we put in the cells "diamonds" \diamond , $V_1^{n,i}$ constitute the third one, and so on. The first column designates the order of singularity of the generalized solution.

Theorem 1.1 asserts that the generalized solution of Problem P1 is bounded if the right-hand side f is orthogonal to all the functions $V_k^{n,i}$ from the table. Theorem 1.2 specifies that (if $l \geq 2$) the singularity of the generalized solution is no worse than r^{1-s} if f is orthogonal to $V_k^{n,i}$ from the triangle under the (s-1)-st row. In other words, the functions from the k-th row of the table are "responsible" for the generalized solutions with behavior r^{-k} near the origin O.

The present paper is a generalization, extension and improvement of the results obtained in [26]. It consists of an introduction and five consecutive sections. Section 2 is devoted to the solutions of the homogeneous adjoint Problems $P1^*$ and $P2^*$. In Section 3 are formulated the 2-D boundary Problems P12 and P13, shortly related to the 3-D Problem P1. The main technical results are established in Sections 4 and 5 – we study the behavior of solutions of 2-D Problems P12 and P13. In the last Section 6 we give proofs of Theorem 1.1 and Theorem 1.2 based on the results of the previous two sections.

2. PROPERTIES OF THE SOLUTIONS H_K^N AND E_K^N

First, we will present three different ways to introduce the solutions of the homogeneous adjoint Problems $P1^*$ and $P2^*$. The functions H_k^n and E_k^n could be found in Khe [20] in the form

$$t\rho^{n-2k-3}(1-t^2/\rho^2)^{n-2k-3/2}F(n-k,-k;3/2;t^2/\rho^2)$$

and

$$\rho^{n-2k-1}(1-t^2/\rho^2)^{n-2k-1/2}F(n-k,-k;1/2;t^2/\rho^2),$$

where F is the hypergeometric Gaus function.

On the other hand, one could obtain $H_k^n(\rho,t)$ and $E_k^n(\rho,t)$ by differentiation of $E_0^n(\rho,t)$ with respect to t.

Lemma 2.1. (see [13, Theorem 4.2]) The functions $H_k^n(\rho,t)$ and $E_k^n(\rho,t)$, defined in Lemma 1.1, satisfy

$$\begin{split} \frac{\partial}{\partial t} H_k^n(\rho,t) &= 2(n-k-1) E_{k+1}^n(\rho,t), \\ \frac{\partial}{\partial t} E_k^n(\rho,t) &= 2(k-n+1/2) H_k^n(\rho,t) \end{split}$$

and they represent some derivative of $E_0^n(\rho,t)$ over t:

$$H_k^n(\rho,t) = \frac{(-1)^{k+1}}{(2n-2k-1)_{2k+1}} \left(\frac{\partial}{\partial t}\right)^{2k+1} \left(\frac{(\rho^2-t^2)^{n-1/2}}{\rho^n}\right),$$

$$E_k^n(\rho, t) = \frac{(-1)^k}{(2n - 2k)_{2k}} \left(\frac{\partial}{\partial t}\right)^{2k} \left(\frac{(\rho^2 - t^2)^{n - 1/2}}{\rho^n}\right).$$

Remark 2.1. This procedure of differentiating the function E_0^n (or $W_0^{n,i}$) with respect to t will (as in Lemma 1.1) produce solutions $V_k^{n,i}$ and $W_k^{n,i}$ of the equation (1.1), but for $k \ge n/2$ the smoothness at the point O will be lost.

Remark 2.2. The solutions of the adjoint Problem $P2^*$ ($P1^*$) given in Lemma 1.1 are not orthogonal. For example, one could check this out for $W_0^{n,i}$ and $W_1^{n,i}$. It is sufficient to show that

$$K := \int_{0}^{1/2} \int_{t}^{1-t} E_0^n(\rho, t) E_1^n(\rho, t) \rho d\rho dt \neq 0.$$

In fact, Lemma 2.1 implies $\partial^2 E_0^n/\partial t^2(\rho,t)=cE_1^n(\rho,t)$ for some constant c and therefore

$$\begin{split} cK &= \int\limits_0^{1/2} \int\limits_t^{1-t} E_0^n(\rho,t) \frac{\partial^2 E_0^n}{\partial t^2}(\rho,t) \rho \; d\rho dt \\ &= \int\limits_{1/2}^1 E_0^n(\rho,1-\rho) \frac{\partial E_0^n}{\partial t} \left(\rho,1-\rho\right) \rho \; d\rho - \int\limits_0^{1/2} \int\limits_t^{1-t} \left(\frac{\partial E_0^n}{\partial t}(\rho,t)\right)^2 \rho \; d\rho dt < 0, \end{split}$$

because $E_0^n(\rho, t) \ge 0$ and $\partial E_0^n/\partial t(\rho, t) \le 0$ for $t \le \rho$.

Remark 2.3. The functions $H_k^n(\rho, t)$ and $E_k^n(\rho, t)$ are linearly independent. Indeed, suppose that some linear combination of these functions is zero. Then from Lemma 2.1 it follows that E_0^n as a function of t is a solution (for a fixed ρ) of a homogeneous linear differential equation with constant coefficients. Therefore E_0^n must be a finite sum of quasi-polynomials of t, which obviously is not true.

A basic tool for our treatment of Protter problems are the Legendre functions P_{ν} (see (3.8) below). Some properties of the Legendre functions P_{ν} one can find in [9]. The next lemma plays a key role in the last section.

Lemma 2.2. Let us denote

$$h_k^{\nu}(\xi,\eta) = \int_{\eta}^{\xi} s^k P_{\nu} \left(\frac{\xi \eta + s^2}{s(\xi + \eta)} \right) ds.$$

If
$$\nu = n - 1/2$$
, then it hold: (a) for $i = 0, 1, ..., \left[\frac{\nu - 1}{2}\right]$

$$h^{\nu}_{\nu - 2i - 2}(\xi, \eta) \Big|_{\xi = \frac{\rho + t}{2}; \eta = \frac{\rho - t}{2}} = c^{n}_{i} \rho^{1/2} H^{n}_{i}(\rho, t)$$

and (**b**) for $i = 0, 1, ..., \left[\frac{\nu}{2}\right]$

$$\left(\frac{\partial}{\partial \xi} - \frac{\partial}{\partial \eta}\right) h_{\nu-2i}^{\nu}(\xi, \eta) \Big|_{\xi = (\rho+t)/2; \eta = (\rho-t)/2} = d_i^n \rho^{1/2} E_i^n(\rho, t),$$

where the constants c_i^n , $d_i^n \neq 0$.

Proof. (a) We will calculate the integrals $h_k^{\nu}(\xi, \eta)$ using the Mellin transform, given by

$$f^*(s) = \int_{0}^{\infty} t^{s-1} f(t) dt,$$

and the Mellin convolution "o", defined by

$$(f \circ g)(x) = \int_{0}^{\infty} f\left(\frac{x}{t}\right) g(t) \frac{dt}{t}.$$

Recall that the relation between both is (see [22, formula (1.2)])

$$(f \circ g)^*(s) = f^*(s)g^*(s).$$
 (2.1)

To apply the last formula to $h_k^{\nu}(\xi,\eta)$, let us introduce new variables x,y and z defined by

$$2\sqrt{x} = \sqrt{\frac{\xi}{\eta}} + \sqrt{\frac{\eta}{\xi}}; \quad 2\sqrt{y} = \frac{\sqrt{\xi\eta}}{s} + \frac{s}{\sqrt{\eta\xi}}; \quad z = \sqrt{\xi\eta}.$$

Then we have

$$\begin{split} \frac{\xi \eta + s^2}{s(\xi + \eta)} &= \sqrt{\frac{y}{x}} \; ; \; \frac{s}{\sqrt{\eta \xi}} = \sqrt{y} \pm \sqrt{y - 1}; \\ \frac{d(2\sqrt{y})}{ds} &= \frac{d}{ds} \left(\frac{\sqrt{\xi \eta}}{s} + \frac{s}{\sqrt{\eta \xi}} \right) = \frac{s^2 - z^2}{s^2 z}; \\ (\sqrt{y} + \sqrt{y - 1})^2 - 1 &= 2\sqrt{y - 1}(\sqrt{y - 1} + \sqrt{y}); \\ (\sqrt{y} - \sqrt{y - 1})^2 - 1 &= 2\sqrt{y - 1}(\sqrt{y - 1} - \sqrt{y}); \end{split}$$

when $s=\sqrt{\xi\eta}$, we have y=1, and y=x for $s=\xi$ or η . Substituting in h_k^{ν} , we find

$$h_k^{\nu}(x,z) = \frac{z^{k+1}}{2} \int_0^{\infty} \left(\frac{x}{y} - 1\right)_+^0 P_{\nu} \left(\sqrt{\frac{y}{x}}\right) (y - 1)_+^{-1/2} \left((\sqrt{y} + \sqrt{y - 1})^{k+1} + (\sqrt{y} - \sqrt{y - 1})^{k+1} \sqrt{y} \frac{dy}{y}\right).$$

Now we are ready to use (2.1) and formulae (11.13(4)), (2.10(4)) from [22] – applying the Mellin transform over x (here " \mapsto " means "transforms into"):

$$(x-1)^0_+ P_{\nu} \left(\sqrt{\frac{1}{x}} \right) \mapsto \frac{\Gamma(-s)\Gamma\left(\frac{1}{2} - s\right)}{\Gamma\left(\frac{1-\nu}{2} - s\right)\Gamma\left(1 + \frac{\nu}{2} - s\right)};$$

$$\sqrt{y}(y-1)_{+}^{-1/2} \left((\sqrt{y} + \sqrt{y-1})^{k+1} + (\sqrt{y} - \sqrt{y-1})^{k+1} \right)$$

$$\mapsto 2\pi \frac{\Gamma\left(\frac{1+k+1}{2} - s - \frac{1}{2} \right) \Gamma\left(\frac{1-k-1}{2} - s - \frac{1}{2} \right)}{\Gamma\left(1 - s - \frac{1}{2} \right) \Gamma\left(\frac{1}{2} - s - \frac{1}{2} \right)}$$

and therefore

$$h_k^{\nu}(x,z)z^{-k-1} \mapsto \pi \frac{\Gamma\left(\frac{1+k}{2} - s\right)\Gamma\left(-\frac{1+k}{2} - s\right)}{\Gamma\left(\frac{1-\nu}{2} - s\right)\Gamma\left(1 + \frac{\nu}{2} - s\right)}.$$
 (2.2)

To find the inverse image of the right-hand side, denote as usual by I_{0+}^{α} the Riemann-Liouville fractional integral (derivative) of order α :

$$\begin{split} I^{\alpha}_{0+}f(x) &:= \frac{1}{\Gamma(\alpha)} \int\limits_{0}^{x} f(t)(x-t)^{\alpha-1} dt, \ \alpha > 0; \\ I^{0}_{0+}f(x) &:= f(x); \\ I^{-\alpha}_{0+}f(x) &:= \left(I^{\alpha}_{0+}\right)^{-1} f(x) := \left(\frac{d}{dx}\right)^{[\alpha]+1} \left(I^{1-\{\alpha\}}_{0+}f\right)(x) \ , \ \alpha > 0. \end{split}$$

Under these notations, the right-hand side of (2.2) is the Mellin transform of the function

$$\frac{\pi}{\Gamma\left(\frac{\nu-k+1}{2}\right)} x^{\frac{\nu+1}{2}} I_{0+}^{\frac{k-\nu+2}{2}} \left(x^{-\frac{k+\nu+3}{2}} (x-1)_{+}^{\frac{\nu-k-1}{2}}\right).$$

Indeed, using (2.2.(4)), (1.10) and (1.4) from [22], for $k < \nu - 1$ we find

$$x^{\frac{\nu+1}{2}} I_{0+}^{\frac{k-\nu+2}{2}} \left(x^{-\frac{k+\nu+3}{2}} (x-1)_{+}^{\frac{\nu-k-1}{2}} \right) \mapsto \Gamma \left(\frac{\nu-k+1}{2} \right) \frac{\Gamma \left(-\frac{1+k}{2} - s \right) \Gamma \left(\frac{1+k}{2} - s \right)}{\Gamma \left(\frac{1-\nu}{2} - s \right) \Gamma \left(1 + \frac{\nu}{2} - s \right)}.$$

Hence, taking $k = \nu - 2i - 2$, for some constants $C_{\nu,i}$ we have

$$h_{\nu-2i-2}^{\nu}(x,z)z^{-\nu+2i+1} = C_{\nu,i}x^{\frac{\nu+1}{2}} \left(\frac{d}{dx}\right)^{i} \left(x^{-\nu+i-1/2}(x-1)_{+}^{i+1/2}\right).$$

Let us now return to the variables ρ and t:

$$\begin{split} z^{\nu-2i-1}x^{\frac{\nu+1}{2}} & \quad (d/dx)^i \left(x^{-\nu+i-1/2}(x-1)^{i+1/2}\right) \\ & = (\xi\eta)^{\frac{\nu-2i-1}{2}} \sum_{j=0}^i d_j' \left(\frac{(\xi+\eta)^2}{\xi\eta}\right)^{-\nu/2+j} \left(\frac{(\xi-\eta)^2}{\xi\eta}\right)^{i+1/2-j} \\ & = (\rho^2-t^2)^{\frac{\nu-2i-1}{2}} \sum_j d_j'' \left(\frac{\rho^2}{(\rho^2-t^2)}\right)^{-\nu/2+j} \left(\frac{t^2}{(\rho^2-t^2)}\right)^{i+1/2-j} \\ & = \sum_{j+l \leq i} d_{j,l} t (\rho^2-t^2)^{\nu-i-1-(j+l)} \rho^{-\nu+2(j+l)} \\ & = \sum_{j=0}^i a_j^i \frac{t (\rho^2-t^2)^{n-3/2-j-i}}{\rho^{n-2j-1/2}}. \end{split}$$

In order to determine the coefficients a_j^i , one could notice that from the definition of h_k^{ν} the function

$$\rho^{-1/2}h^{\nu}_{\nu-2i-2}\left(\frac{\rho+t}{2},\frac{\rho-t}{2}\right)\sin n\varphi$$

satisfies the wave equation. Therefore, after substituting in the equation, we find $a_i^i = a_0^i A_i^i$ and

$$\rho^{-1/2}h_{\nu-2i-2}^{\nu}\left(\frac{\rho+t}{2},\frac{\rho-t}{2}\right) = a_i^{\nu}H_i^n(\rho,t).$$

(b) Let us find the functions $(\partial/\partial \xi - \partial/\partial \eta) h^{\nu}_{\nu-2i}(\xi,\eta)$, where $i=0,1,...,\left[\frac{\nu}{2}\right]$. Notice that for $i \geq 1$, due to Lemma 2.1 and (a),

$$\rho^{-1/2} \left(\frac{\partial}{\partial \xi} - \frac{\partial}{\partial \eta} \right) h^{\nu}_{\nu - 2i}(\xi, \eta) = C^n_i \frac{\partial}{\partial t} H^n_{i-1}(\rho, t) = C^{n'}_i E^n_i(\rho, t).$$

Only the case i=0 has been omitted, i.e. we need to calculate

$$\left(\frac{\partial}{\partial \xi} - \frac{\partial}{\partial \eta}\right) h_{\nu}^{\nu}(\xi, \eta).$$

Therefore we consider the function

$$\frac{\partial}{\partial t} z^{k+1} x^{\frac{\nu+1}{2}} I_{0+}^{\frac{k-\nu+2}{2}} \left(x^{-\frac{k+\nu+3}{2}} (x-1)_{+}^{\frac{\nu-k-1}{2}} \right),$$

where $k = \nu$ and

$$x = \frac{\rho^2}{\rho^2 - t^2}; \quad 2z = (\rho^2 - t^2)^{1/2}.$$

That is

$$\begin{split} &\frac{\partial}{\partial t} z^{\nu+1} x^{\frac{\nu+1}{2}} I_{0+}^{1} \left(x^{-\nu - \frac{3}{2}} (x-1)_{+}^{-1/2} \right) = \frac{\partial}{\partial t} z^{\nu+1} x^{\frac{\nu+1}{2}} \int_{0}^{x} \tau^{-\nu - \frac{3}{2}} (\tau-1)_{+}^{-1/2} d\tau \\ &= \frac{\partial}{\partial t} \rho^{\nu+1} 2^{-\nu-1} \int_{0}^{x} \tau^{-\nu - \frac{3}{2}} (\tau-1)^{-1/2} d\tau = 2^{-\nu-1} \rho^{\nu+1} x^{-\nu - \frac{3}{2}} (x-1)^{-1/2} \frac{\partial x}{\partial t} \\ &= 2^{-\nu-1} \rho^{\nu+1} \frac{(\rho^{2} - t^{2})^{\nu+3/2}}{\rho^{2\nu+3}} \frac{(\rho^{2} - t^{2})^{1/2}}{t} \frac{2t\rho^{2}}{(\rho^{2} - t^{2})^{2}} = 2^{-\nu} \frac{(\rho^{2} - t^{2})^{\nu}}{\rho^{\nu}}. \end{split}$$

Hence for $\nu = n - 1/2$ we conclude that

$$\rho^{-1/2} \left(\frac{\partial}{\partial \xi} - \frac{\partial}{\partial \eta} \right) h_{\nu}^{\nu}(\xi, \eta) = \rho^{-1/2} C_n \frac{(\rho^2 - t^2)^{n-1/2}}{\rho^{n-1/2}} = C_n E_0^n(\rho, t). \quad \Box$$

3. SOLVING PROBLEM P1

In terms of Theorems 1.1 and 1.2 it is sufficient to study the Problem P1 only when the right-hand side f of the wave equation is simply

$$f(\rho, t, \varphi) = f_n^1(\rho, t) \cos n\varphi + f_n^2(\rho, t) \sin n\varphi, \quad n \in \mathbb{N} \cup \{0\}.$$

Then we seek solutions of the wave equation of the same form:

$$u(\rho, t, \varphi) = u_n^1(\rho, t) \cos n\varphi + u_n^2(\rho, t) \sin n\varphi.$$

Thus Problem P1 reduces to the following one:

Problem P12. To solve the equation

$$(u_n)_{\rho\rho} + \frac{1}{\rho}(u_n)_{\rho} - (u_n)_{tt} - \frac{n^2}{\rho^2}u_n = f_n \tag{3.1}$$

in $\Omega_0 = \{0 < t < 1/2; t < \rho < 1 - t\} \subset \mathbb{R}^2$ with the boundary conditions

P12 $u_n(\rho, 0) = 0$ for $0 < \rho \le 1$ and $u_n(\rho, 1 - \rho) = 0$ for $1/2 \le \rho \le 1$.

Let us now introduce new coordinates

$$\xi = \frac{\rho + t}{2} \; ; \; \; \eta = \frac{\rho - t}{2},$$
 (3.2)

and set

$$v(\xi, \eta) = \rho^{1/2} u_n(\rho, t) \; ; \; g(\xi, \eta) = \rho^{1/2} f_n(\rho, t).$$
 (3.3)

Denoting $\nu = n - \frac{1}{2}$, one transforms Problem P12 into

Problem P13. To find a solution $v(\xi, \eta)$ of the equation

$$v_{\xi\eta} - \frac{\nu(\nu+1)}{(\xi+\eta)^2}v = g \tag{3.4}$$

in the domain $D = \{0 < \xi < 1/2; 0 < \eta < \xi\}$ with the following boundary conditions:

P13
$$v(\xi,\xi) = 0$$
 for $\xi \in (0,1/2)$ and $v(1/2,\eta) = 0$ for $\eta \in (0,1/2)$.

Problems P12 and P13 have been introduced in [25], although the change of coordinates $\xi = 1 - \rho - t$ and $\eta = 1 - \rho + t$ is used there instead of (3.2). Of course, because the solution of Problem P1 may be singular, the same is true for the solutions of P12 and P13. For that reason, Popivanov and Schneider [25] have defined and proved the existence and uniqueness of generalized solutions of Problems P12 and P13, which correspond to the generalized solution of Problem P1. Further, by "solution" of Problem P12 or P13 we will mean exactly this unique generalized solution.

Remark 3.1. Notice that even when the right-hand side function $f_n(\rho, t)$ belongs to $C^k(\overline{\Omega}_0)$, the corresponding function $g(\xi, \eta) = \rho^{1/2} f_n(\rho, t)$ in (3.4) belongs to $C^k(\overline{\Omega}_0 \setminus O)$, but its derivative may not be continuous at the origin O. At the same time, when the solution $v(\xi, \eta)$ of Problem P13 is bounded, the solution $u_n(\rho, t) = \rho^{-1/2} v(\xi, \eta)$ of Problem P12 may be singular.

Nevertheless, we will solve Problem P13 instead of Problem P1. We can construct the solution of the Problem P13 using two different methods. First, following Popivanov, Schneider [25], one could use the equivalent integral equation

$$U(\xi_0, \eta_0) = \int_{\xi=\xi_0}^{\xi=\frac{1}{2}} \int_{\eta=\eta_0}^{\eta=\xi_0} \left(\frac{\nu(\nu+1)}{(\xi+\eta)^2} U(\xi, \eta) + g(\xi, \eta) \right) d\eta d\xi$$
 (3.5)

and construct the solution as a limit of a sequence of successive approximations $U^{(k)}$ defined by

$$U^{(0)}(\xi_0, \eta_0) = \int_{\xi=\xi_0}^{\xi=\frac{1}{2}} \int_{\eta=\eta_0}^{\eta=\xi_0} g(\xi, \eta) d\eta d\xi,$$
 (3.6)

$$U^{(k+1)}(\xi_0, \eta_0) = \int_{\xi=\xi_0}^{\xi=\frac{1}{2}} \int_{\eta=\eta_0}^{\eta=\xi_0} \left(\frac{\nu(\nu+1)}{(\xi+\eta)^2} U^{(k)}(\xi, \eta) + g(\xi, \eta) \right) d\eta d\xi.$$
 (3.7)

On the other hand, notice that the function

$$R(\xi_1, \eta_1, \xi, \eta) = P_{\nu} \left(\frac{(\xi - \eta)(\xi_1 - \eta_1) + 2\xi_1 \eta_1 + 2\xi \eta}{(\xi_1 + \eta_1)(\xi + \eta)} \right)$$
(3.8)

is a Riemann function for the equation (3.4) (Copson [7]). Therefore, we can construct the function v as a solution of a Goursat's problem in D with boundary conditions $v(1/2, \eta) = 0$ and $v(\xi, 0) = \varphi(\xi)$ with some unknown function $\varphi(\xi)$, which will be determined later:

$$v(\xi,\eta) = \varphi(\xi) + \int_{\xi}^{\frac{1}{2}} \varphi(\xi_1) \frac{\partial}{\partial \xi_1} R(\xi_1,0,\xi,\eta) d\xi_1 - \int_{\xi}^{\frac{1}{2}} \int_{0}^{\eta} R(\xi_1,\eta_1,\xi,\eta) g(\xi_1,\eta_1) d\eta_1 d\xi_1.$$
(3.9)

Now, following Aldashev [1], the boundary condition $v(\xi, \xi) = 0$ gives the equation

$$G(\xi) = \varphi(\xi) + \int_{\xi}^{\frac{1}{2}} \varphi(\xi_1) \frac{\partial}{\partial \xi_1} P_{\nu} \left(\frac{\xi}{\xi_1}\right) d\xi_1 = -\int_{\xi}^{\frac{1}{2}} \varphi'(\xi_1) P_{\nu} \left(\frac{\xi}{\xi_1}\right) d\xi_1$$
 (3.10)

for the function $\varphi(\xi)$, where

$$G(\xi) = \int_{\xi}^{\frac{1}{2}} \int_{0}^{\xi} P_{\nu} \left(\frac{\xi_{1}\eta_{1} + \xi^{2}}{\xi(\xi_{1} + \eta_{1})} \right) g(\xi_{1}, \eta_{1}) d\eta_{1} d\xi_{1}.$$
 (3.11)

According to formulae (35.17), (35.28) from [28], the integral equation (3.10) is invertible and we have

$$\frac{d}{d\xi}\varphi(\xi) = -\xi \frac{d^2}{d\xi^2} \xi \int_{\xi}^{\frac{1}{2}} P_{\nu} \left(\frac{\xi_1}{\xi}\right) \frac{G(\xi_1)}{\xi_1^2} d\xi_1 = G'(\xi) + \frac{d}{d\xi} \int_{\xi}^{\frac{1}{2}} P'_{\nu} \left(\frac{\xi_1}{\xi}\right) \frac{G(\xi_1)}{\xi_1} d\xi_1.$$

Finally, due to $\varphi(1/2) = G(1/2) = 0$, it follows

$$\varphi(\xi) = G(\xi) + \int_{\xi}^{\frac{1}{2}} P_{\nu}' \left(\frac{\xi_1}{\xi}\right) \frac{G(\xi_1)}{\xi_1} d\xi_1. \qquad \Box$$
 (3.12)

Using the sequence $\{U^{(k)}\}\$ defined by (3.6), (3.7), we have the following

Lemma 4.1. Let $0 < \varepsilon < 1$ and suppose that $|\nu(\nu+1)| < \varepsilon(1+\varepsilon)$. Then the solution $U(\xi,\eta)$ of the integral equation (3.5) in D satisfies the estimate

$$|U(\xi,\eta)| \le C(\xi-\eta)(\xi+\eta)^{-\varepsilon} \sup_{D} |g|,$$

where the constant C depends only on ν and ε .

Proof. The key point in the proof is the estimate for the integrals:

$$\begin{split} I_{\varepsilon} &:= \int\limits_{\xi=\xi_{0}}^{\xi=\frac{1}{2}} \int\limits_{\eta=\eta_{0}}^{\eta=\xi_{0}} (\xi-\eta)(\xi+\eta)^{-\varepsilon-2} d\eta d\xi \\ &= \int\limits_{\xi=\frac{1}{2}}^{\xi=\frac{1}{2}} \int\limits_{\eta=\xi_{0}}^{\eta=\xi_{0}} \left((\xi+\eta)^{-\varepsilon-1} - 2\eta(\xi+\eta)^{-\varepsilon-2} \right) d\eta d\xi \\ &= \int\limits_{\eta_{0}}^{\xi} \left(\frac{1}{\varepsilon} (\xi_{0}+\eta)^{-\varepsilon} - \frac{1}{\varepsilon} (\frac{1}{2}+\eta)^{-\varepsilon} - \frac{2\eta}{1+\varepsilon} (\xi_{0}+\eta)^{-\varepsilon-1} + \frac{2\eta}{1+\varepsilon} (\frac{1}{2}+\eta)^{-\varepsilon-1} \right) d\eta \\ &\leq \int\limits_{\eta_{0}}^{\xi_{0}} \left(\frac{1}{\varepsilon} (\xi_{0}+\eta)^{-\varepsilon} - \frac{2\eta}{1+\varepsilon} (\xi_{0}+\eta)^{-\varepsilon-1} \right) d\eta \\ &= \int\limits_{\eta_{0}}^{\xi_{0}} \left(\frac{1-\varepsilon}{\varepsilon(1+\varepsilon)} (\xi_{0}+\eta)^{-\varepsilon} + \frac{2\xi_{0}}{1+\varepsilon} (\xi_{0}+\eta)^{-\varepsilon-1} \right) d\eta \\ &= -\frac{1}{\varepsilon(1+\varepsilon)} (\xi_{0}+\eta_{0})^{1-\varepsilon} + \frac{2\xi_{0}}{\varepsilon(1+\varepsilon)} (\xi_{0}+\eta_{0})^{-\varepsilon} \\ &= \frac{1}{\varepsilon(1+\varepsilon)} (\xi_{0}-\eta_{0})(\xi_{0}+\eta_{0})^{-\varepsilon}. \end{split}$$

For the function $U^{(0)}$ we find

$$\sup_{\Omega} |U^{(0)}| \le \frac{1}{2} (\xi_0 - \eta_0) \sup_{D} |g| \le (\xi_0 - \eta_0) (\xi_0 + \eta_0)^{-\varepsilon} \sup_{D} |g|.$$

Now, because of

$$\left| (U^{(k+1)} - U^{(k)})(\xi_0, \eta_0) \right| \leq \int_{\xi = \xi_0}^{\xi = \frac{1}{2}} \int_{\eta = \eta_0}^{\eta = \xi_0} \frac{|\nu(\nu + 1)|}{(\xi + \eta)^2} \left| (U^{(k)} - U^{(k-1)})(\xi, \eta) \right| d\eta d\xi,$$

if take
$$\alpha := \frac{|\nu(\nu+1)|}{\varepsilon(1+\varepsilon)}$$
, we have by induction
$$\left|(U^{(k)} - U^{(k-1)})(\xi_0, \eta_0)\right| \le \alpha^k (\xi_0 - \eta_0)(\xi_0 + \eta_0)^{-\varepsilon} \sup_{\Omega} |g|.$$

Indeed, using the calculations of I_{ε} , we find that

$$\left| (U^{(k+1)} - U^{(k)})(\xi_0, \eta_0) \right| \leq \int_{\xi = \xi_0}^{\xi = \frac{1}{2}} \int_{\eta = \eta_0}^{\eta = \xi_0} \frac{|\nu(\nu + 1)|}{(\xi + \eta)^{2+\varepsilon}} \alpha^k (\xi - \eta) \sup_{D} |g| d\eta d\xi
\leq \alpha^{k+1} (\xi_0 - \eta_0) (\xi_0 + \eta_0)^{-\varepsilon} \sup_{D} |g|.$$
(4.1)

Finally, we arrive at the estimate

$$|U^{(k+1)}(\xi_{0}, \eta_{0})| \leq \sum_{i=0}^{k} \left| (U^{(i+1)} - U^{(i)})(\xi_{0}, \eta_{0}) \right|$$

$$\leq \sum_{i=0}^{k} \alpha^{i+1} (\xi_{0} - \eta_{0})(\xi_{0} + \eta_{0})^{-\varepsilon} \sup_{D} |g|$$

$$= \frac{(1 - \alpha^{k+2})}{(1 - \alpha)} (\xi_{0} - \eta_{0})(\xi_{0} + \eta_{0})^{-\varepsilon} \sup_{D} |g|,$$

which, together with (4.1), shows that if $\alpha < 1$, the sequence $\{U^{(k)}\}$ uniformly converges to a solution $U(\xi, \eta)$ of

$$U(\xi_0, \eta_0) = \int_{\xi=\xi_0}^{\xi=\frac{1}{2}} \int_{\eta=\eta_0}^{\eta=\xi_0} \left(\frac{\nu(\nu+1)}{(\xi+\eta)^2} U(\xi, \eta) + g(\xi, \eta) \right) d\eta d\xi.$$

Even more, it holds the estimate

$$|U(\xi,\eta)| \le \frac{1}{(1-\alpha)} (\xi-\eta)(\xi+\eta)^{-\varepsilon} \sup_{D} |g|.$$

Now, for n=0 or 1 we have $\nu=-1/2$ or 1/2, respectively, and we will apply Lemma 4.1 with suitable ε .

Theorem 4.1. For the solution $u_n(\rho,t)$ of Problem P12 with right-hand side function $f_n \in C(\overline{\Omega}_0)$ the following a priori estimate holds:

(i) for the case n = 0

$$|u_0(\rho,t)| \le C\rho^{1/4} \max_{\overline{\Omega}_0} |\rho^{1/2} f_0|$$

with constant C, independent of the function $f_0(\rho, t)$;

(ii) if n = 1, then for every $\delta > 0$ there exists a constant C_{δ} , independent of the function $f_1(\rho, t)$, such that

$$|u_1(\rho,t)| \le C_\delta \rho^{-\delta} \max_{\overline{\Omega}_0} |\rho^{1/2} f_1|.$$

Proof. Notice that when n=0 and n=1 we have $|\nu(\nu+1)|=1/4$ and $|\nu(\nu+1)|=3/4$, respectively. Therefore one could apply Lemma 4.1 for the solution

 $v(\xi,\eta) = U(\xi,\eta)$ of Problem P13 with $\varepsilon = 1/4$ and $\varepsilon = 1/2 + \delta$, respectively. The assertion follows from the relation (3.3):

$$|u(\rho,t)| = \left|\rho^{-1/2}v\left(\frac{\rho+t}{2},\frac{\rho-t}{2}\right)\right| \leq C'\rho^{1-\varepsilon-1/2}\max_{\overline{D}}|g| \leq C\rho^{1/2-\varepsilon}\max_{\overline{\Omega}_0}|\rho^{1/2}f|.\,\square$$

5. THE CASE
$$n > 2$$

In this case, unlike the above approach, the behavior of the solution is studied in [26], using the properties of the Riemann function (3.8), given by Legendre functions P_{ν} . Let us remind some of these results here. For the function $v(\xi,\eta)$, defined by (3.9), (3.12) and (3.11), it is not hard to see that $v(\xi,0) = \varphi(\xi)$ may blow up when ξ tends to 0. Nevertheless, one could control the growth of

$$I(\xi) := \int_{\xi}^{\frac{1}{2}} P_{\nu}' \left(\frac{\xi_1}{\xi}\right) \frac{G(\xi_1)}{\xi_1} d\xi_1 \tag{5.1}$$

with the help of the following lemma:

Lemma 5.1. (see [26, Lemma 3.1]) Let $\nu > 1$ be a real number with integral part $[\nu]$ and fractional part $\{\nu\} = \nu - [\nu] \neq 0$. Suppose that $G \in C^{[\nu-2]+}(0,1/2]$, $|G^{(k)}(\xi)| \leq A\xi^{1-k}$ for $k = 0, 1, ..., [\nu - 2]_+$ for some constant A, and

$$\int_{0}^{\frac{1}{2}} \xi^{\nu - 2i - 2} G(\xi) d\xi = 0 \quad \text{for } i = 0, 1, ..., \left[\frac{\nu - 1}{2} \right]. \tag{5.2}$$

Then the function $I(\xi)$, defined by (5.1), is $C^{[\nu]-1}(0,1/2]$. More precisely, there is a constant C, independent of $G(\xi)$, such that:

- (i) |I(ξ)| ≤ CAξ if [ν] is an odd number;
 (ii) |I(ξ)| ≤ CAξ^{1-{ν}} if [ν] is an even number.

Besides, Lemma 3.2 from [26] asserts that each of the orthogonality conditions (5.2) actually "controls" one power-type singularity of the function $I(\xi)$:

Lemma 5.2. Let $\nu > 1, \{\nu\} \neq 0$, and p be a nonnegative integer, $p \leq$ $[(\nu-1)/2]$. Suppose that $G \in C^{(2p-1)+}(0,1/2], |G^{(k)}(\xi)| \leq A\xi^{2p-[\nu]+1-k}$ for k=1 $0, 1, ..., (2p-1)_+$ for some constant A, and

$$\int_{0}^{\frac{1}{2}} \xi^{\nu-2i-2} G(\xi) d\xi = 0 \quad \text{for } i = 0, 1, ..., p-1.$$
 (5.3)

Then the estimate

$$|I(\xi)| \le C_1 A \xi^{-(\nu - 2p - 1)}$$

holds for some constant C_1 , independent of $G(\xi)$. Moreover, if

$$\int_{0}^{\frac{1}{2}} \xi^{\nu-2p-2} G(\xi) d\xi \neq 0, \tag{5.4}$$

then

$$|I(\xi)| \ge C_2 \xi^{-(\nu - 2p - 1)}$$

for $C_2 > 0$ and sufficiently small ξ .

In other words, one has to impose some conditions on the right-hand side g of Problem P13 to secure certain behavior of the solution v. In fact, the definition (3.11) of the function $G(\xi)$ gives the equality (see [26])

$$\int_{0}^{\frac{1}{2}} \xi^{\nu-2i-2} G(\xi) d\xi = \int_{0}^{\frac{1}{2}} \int_{0}^{\xi_{1}} \left(\int_{\eta_{1}}^{\xi_{1}} \xi^{\nu-2i-2} P_{\nu} \left(\frac{\xi_{1}\eta_{1} + \xi^{2}}{\xi(\xi_{1} + \eta_{1})} \right) d\xi \right) g(\xi_{1}, \eta_{1}) d\eta_{1} d\xi_{1}.$$

It shows that one needs orthogonality of g to the functions $h_{\nu-2i-2}^{\nu}$ defined in Lemma 2.2. As a result we are able to prove the following

Theorem 5.1. (see [26, Theorem 4.1]) Let $n \in \mathbb{N}$, $n \geq 2$. Suppose that $g \in C^{(n-4)+}(\overline{D}\backslash O)$, $|D^{\alpha}g(\xi,\eta)| \leq A\xi^{-|\alpha|}$ for $|\alpha| \leq (n-4)_+$, and the orthogonality conditions

$$\int_{0}^{\frac{1}{2}} \int_{0}^{\xi} h_{\nu-2i-2}^{\nu}(\xi,\eta) g(\xi,\eta) d\eta d\xi = 0 \quad \text{for } i = 0, 1, ..., \left[\frac{\nu-1}{2} \right]$$
 (5.5)

are satisfied with $\nu = n - 1/2$. Then the solution $v(\xi, \eta)$ of Problem P13 belongs to $C^{(n-4)+1}(\overline{D}\backslash O)$ and satisfies the following estimates:

(i) if n is an even number, then

$$|v(\xi,\eta)| \le CA\xi |\ln \xi|;$$

(ii) if n is an odd number, then

$$|v(\xi,\eta)| \le CA\xi^{1/2}.$$

In both cases the constant C does not depend on the function $g(\xi, \eta)$.

In the same way one gets the following theorem, which corresponds to Lemma 5.2:

Theorem 5.2. (see [26, Theorem 4.2]) Let p be a nonnegative integer and $p \leq [(\nu-1)/2]$. Suppose that the function $g \in C^{(2p-2)+}(\overline{D}\backslash O)$, $|D^{\alpha}g(\xi,\eta)| \leq A\xi^{-|\alpha|}$ for $|\alpha| \leq (2p-2)_+$ and

$$\int_{0}^{\frac{1}{2}} \int_{0}^{\xi} h_{\nu-2i-2}^{\nu}(\xi, \eta) g(\xi, \eta) d\eta d\xi = 0 \quad \text{for } i = 0, 1, ..., p-1.$$
 (5.6)

Then the estimate

$$|v(\xi,\eta)| \le CA\xi^{-(\nu-2p-1)}$$
 (5.7)

holds for some constant C, independent of the function $g(\xi,\eta)$. Moreover, the condition

$$\int_{0}^{\frac{1}{2}} \int_{0}^{\xi} h_{\nu-2p-2}^{\nu}(\xi, \eta) g(\xi, \eta) d\eta d\xi \neq 0$$
 (5.8)

implies that the lower estimate

$$|v(\xi,0)| \ge c\xi^{-(\nu-2p-1)} \tag{5.9}$$

holds for some constant c > 0 and sufficiently small ξ .

All these preparations and Lemma 2.2 lead us to the following estimate for the solution of Problem P12:

Theorem 5.3. Suppose that $n \in \mathbb{N}$, $n \geq 2$, $f_n(\rho,t) \in C^{(n-4)+}(\overline{\Omega}_0 \setminus O)$, and there is a constant A such that $|\rho^{|\alpha|+1/2}D^{\alpha}f_n(\rho,t)| \leq A$ for $|\alpha| \leq (n-4)_+$. Let also there hold the orthogonality conditions

$$\int_{0}^{\frac{1}{2}} \int_{t}^{1-t} H_{i}^{n}(\rho, t) f_{n}(\rho, t) \rho d\rho dt = 0 \quad \text{for } i = 0, 1, ..., \left[\frac{n}{2}\right] - 1.$$
 (5.10)

Then the solution $u_n(\rho,t)$ of Problem P12 satisfies:

- (i) $|u_n(\rho,t)| \le CA\rho^{1/2} |\ln \rho|$ if n is an even number;
- (ii) $|u_n(\rho, t)| \leq CA$ if n is an odd number.

In both cases the constant C does not depend on the function $f_n(\rho, t)$.

Proof. Let us define the function $g(\xi,\eta) = \rho^{1/2} f_n(\rho,t)$, where $\xi = (\rho+t)/2$, $\eta = (\rho-t)/2$. Then the estimates for f_n imply that g satisfies $|D^{\alpha}g(\xi,\eta)| \leq CA\xi^{-|\alpha|}$ for $|\alpha| \leq (n-4)_+$. The orthogonality conditions, due to Lemma 2.2, yield

$$\int\limits_{0}^{\frac{1}{2}}\int\limits_{0}^{\xi}h_{\nu-2i-2}^{\nu}(\xi,\eta)g(\xi,\eta)\;d\eta d\xi\;=C\int\limits_{0}^{\frac{1}{2}}\int\limits_{t}^{1-t}H_{i}^{n}(\rho,t)f_{n}(\rho,t)\rho d\rho dt=0.$$

Now Theorem 5.1 gives the required estimates for the solution $v(\xi, \eta)$ of Problem P13, given by (3.9), and $u = \rho^{-1/2}v$ is the solution of Problem P12.

Using the similar arguments, we get the corresponding result for the case when not all of the orthogonality conditions (5.10) are fulfilled:

Theorem 5.4. Let $n, q \in \mathbb{N} \cup \{0\}$, $n \geq 2$, $q \leq \left[\frac{n}{2}\right] - 1$. Suppose that the function $f_n \in C^{(2q-2)+}(\overline{\Omega}_0 \backslash O)$, $|D^{\alpha}f_n(\rho,t)| \leq A\rho^{-|\alpha|-1/2}$ for $|\alpha| \leq (2q-2)_+$ and

$$\int_{0}^{\frac{1}{2}} \int_{t}^{1-t} H_{i}^{n}(\rho, t) f_{n}(\rho, t) \rho d\rho dt = 0 \quad \text{for } i = 0, 1, ..., q - 1.$$
 (5.11)

Then for the solution $u_n(\rho,t)$ of Problem P12 the upper estimate

$$|u_n(\rho, t)| \le CA\rho^{-(n-2q-1)}$$
 (5.12)

holds, where the constant C is independent of $f_n(\rho, t)$. If we suppose also that

$$\int_{0}^{\frac{1}{2}} \int_{t}^{1-t} H_{q}^{n}(\rho, t) f_{n}(\rho, t) \rho d\rho dt \neq 0,$$
(5.13)

then the lower estimate

$$|u_n(\rho,\rho)| \ge c\rho^{-(n-2q-1)}$$

holds for c > 0 and sufficiently small ρ .

Proof. We again define $g(\xi,\eta) = \rho^{1/2} f_n(\rho,t)$ and we prove the theorem applying Theorem 5.2 for $v = \rho^{1/2} u$ and g instead of Theorem 5.1 as in the proof of Theorem 5.3.

Finally, Theorems 5.3 and 5.4 show that every solution of Problem P12 is a linear combination of at most $\left[\frac{n}{2}\right]$ fixed singular solutions:

Lemma 5.3. For $n \geq 2$ there exist $\left[\frac{n}{2}\right]$ functions $v_n^i(\rho, t)$, $i = 0, ..., \left[\frac{n}{2}\right] - 1$, such that for every generalized solution $u_n(\rho, t)$ of Problem P12 with some right-hand side function $f_n \in C^{(n-4)+}(\overline{\Omega}_0)$ the equality

$$u_n(\rho, t) = \sum_{i=0}^{\left[\frac{n}{2}\right]-1} c_i v_n^i(\rho, t) + w(\rho, t)$$

holds with some constants c_i and some bounded function $w(\rho, t)$ dependent on $u_n(\rho, t)$.

Proof. Let $u_n(\rho,t)$ be the generalized solution of Problem P12 with some right-hand side function $f_n(\rho,t) \in C^{(n-4)+}(\overline{\Omega}_0)$. In general, u_n has a singularity at the origin O. Let k be an integer, $0 \le k \le \left[\frac{n}{2}\right] - 1$, and $f_n^{(k)}(\rho,t)$ be the projection of f_n on the linear space $L_{k,n}$ of functions, orthogonal to the functions H_i^n for i = 0, 1, ..., k:

$$L_{k,n} := \left\{ f(\rho,t) : \int_{0}^{\frac{1}{2}} \int_{t}^{1-t} H_{i}^{n}(\rho,t) f(\rho,t) \rho d\rho dt = 0 \text{ for } i = 0, 1, ..., k \right\}.$$

Then $f_n - \sum_{i=0}^k \alpha_i^k H_i^n = f_n^{(k)} \in L_{k,n}$ with some constants α_i^k such that

$$\int_{0}^{\frac{1}{2}} \int_{t}^{1-t} H_{j}^{n}(\rho, t) f_{n}^{(k)}(\rho, t) \rho d\rho dt = 0, \quad j = 0, ..., k,$$

i.e.

$$\sum_{i=0}^k \alpha_i^k \int\limits_0^{\frac12} \int\limits_t^{1-t} H_j^n(\rho,t) H_i^n(\rho,t) \rho d\rho dt = \int\limits_0^{\frac12} \int\limits_t^{1-t} H_j^n(\rho,t) f_n(\rho,t) \rho d\rho dt$$

for j=0,1,...,k. This system has an unique solution for constants α_i^k . To show this, suppose that the rank of the $(k\times k)$ -matrix with elements $\int\limits_0^{\frac{1}{2}}\int\limits_t^{1-t}H_j^nH_i^n\rho d\rho dt$ is less than k. Then there are numbers $\beta_0,\,\beta_1,...,\beta_k$ such that at least one is not zero and

$$\sum_{i=0}^{k} \beta_{i} \int_{0}^{\frac{1}{2}} \int_{t}^{1-t} H_{j}^{n}(\rho, t) H_{i}^{n}(\rho, t) \rho d\rho dt = 0 \quad \text{for } j = 0, 1, ..., k$$

or

$$\int_{0}^{\frac{1}{2}} \int_{t}^{1-t} H_{j}^{n}(\rho, t) \left(\sum_{i=0}^{k} \beta_{i} H_{i}^{n}(\rho, t) \right) \rho d\rho dt = 0 \quad \text{for } j = 0, 1, ..., k.$$

Therefore $\sum_{i=0}^{k} \beta_i H_i^n = 0$ in Ω_0 , which is impossible, because the functions H_i^n are linearly independent in view of Remark 2.3.

Let us denote by $v_n^i(\rho, t)$ the solution of the Problem P12 with a right-hand side function $H_i^n(\rho, t)$, and by $u_n^{(k)}(\rho, t)$ the solution with a right-hand side function $f_n^{(k)}(\rho, t)$. Then we have for u_n the representation

$$u_n = \sum_{i=0}^k \alpha_i^k v_n^i + u_n^{(k)}.$$

When $k < \left[\frac{n}{2}\right] - 1$, Theorem 5.4 gives the estimate

$$|u_n^{(k)}(\rho,t)| \le C\rho^{-(n-2k-3)},$$

while for $k = k_0 := \left[\frac{n}{2}\right] - 1$ Theorem 5.3 shows that the function $w := u_n^{(k_0)}$ is at least bounded.

6. PROOF OF THE MAIN RESULTS

We are now ready to prove Theorem 1.1 and Theorem 1.2 (see Introduction). **Proof of Theorem 1.1.** When the right-hand side function has the form

$$f = f_0^1(\rho, t) + \sum_{n=1}^l \left(f_n^1(\rho, t) \cos n\varphi + f_n^2(\rho, t) \sin n\varphi \right),$$

one could construct the unique generalized solution $u(x_1, x_2, t)$ as

$$u = u_0^1(\rho, t) + \sum_{n=1}^{l} \left(u_n^1(\rho, t) \cos n\varphi + u_n^2(\rho, t) \sin n\varphi \right), \tag{6.1}$$

where the functions u_n^i are solutions of Problem P12 with right-hand side function $f_n^i \in C^{(l-4)+}(\overline{\Omega}_0)$. First of all, we use Theorem 4.1 to estimate the functions u_0^1 , u_1^1 and u_1^2 :

$$|u_0^1(\rho,t)| \le C\rho^{1/4} \max_{\overline{\Omega}_0} |\rho^{1/2} f_0^1| \le Cr^{1/4} ||f_0||_0^*,$$

$$|u_1^i(\rho,t)| \le C_\delta \rho^{-\delta} \max_{\overline{\Omega}_0} |\rho^{1/2} f_1^i| \le C_\delta' r^{-\delta} ||f_1||_0^*.$$

For the case $n \geq 2$ we apply Theorem 5.3 with the constant $A = ||f_n||_{(n-4)_+}^*$. Because of the identity

$$\int_{0}^{\frac{1}{2}} \int_{t}^{1-t} H_{k}^{n}(\rho, t) f_{n}^{i}(\rho, t) \rho d\rho dt = c \int_{\Omega} V_{k}^{n, i}(x_{1}, x_{2}, t) f(x_{1}, x_{2}, t) dx_{1} dx_{2} dt = 0, \quad (6.2)$$

the orthogonality conditions (5.10) are fulfilled. Therefore, if n is an even number,

$$|u_n^i| \le C_n A \rho^{1/2} |\ln \rho| \le C_n' r^{1/2} |\ln r| ||f_n||_{(n-4)_+}^*,$$

while

$$|u_n^i| \le C_n A = C_n ||f_n||_{(n-4)_+}^*$$

if n is an odd number. Finally, summing up all these inequalities, we find

$$|u| \le |u_0^1| + \sum_{n=1}^l (|u_n^1| + |u_n^2|) \le C_1 r^{1/4} ||f_0||_0^* + C_{2,\delta} r^{-\delta} ||f_1||_0^* + C_{3r^{1/2}} |\ln r| \sum_{k=1}^{\lfloor l/2 \rfloor} ||f_{2k}||_{(2k-4)_+}^* + C_4 \sum_{k=1}^{\lfloor \frac{l-1}{2} \rfloor} ||f_{2k+1}||_{(2k-3)_+}^*. \quad \Box$$

Remark 6.1. Notice that the orthogonality condition $\int_{\Omega} V_k^{n,i} f dx_1 dx_2 dt = 0$ for a function f with the representation (1.3) is imposed only on the function $f_n^i(\rho,t)$ and has no influence over the other functions $f_m^j(\rho,t)$ from (1.3) with indices $(m,j) \neq (n,i)$.

Proof of Theorem 1.2. We use again the representation (6.1). For n = 0 and I we have respectively

$$|u_0^1| \le C\rho^{1/4}||f_0||_0^* \le C'r^{-s+1}||f_0||_0^*$$

and

$$|u_1^i| \le C_{1/2} \rho^{-1/2} ||f_1||_0^* \le C'' r^{-s+1} ||f_1||_0^*,$$

because $s \ge 2$. For $n \ge 2$ we apply Theorem 5.4 with $q = \left[\frac{n-s+1}{2}\right]_+$ and $A = ||f_n||_{2q-2}^*$. Now, the identity (6.2) shows that the orthogonality conditions

hold for $n-2k \ge s+1$, i.e. for all $k \ge 0$ such that $q-1 \ge \left[\frac{n-s+1}{2}\right]-1 \ge \left[\frac{2k+2}{2}\right]-1=k$. Therefore

$$|u_n^i| \le CA\rho^{-n+2q+1} \le C'\rho^{-s+1}||f_n||_{(2q-2)_+}^* \le C''r^{-s+1}||f_n||_{(n-s-1)_+}^*,$$

because depending of parity of n we have $-n+2\left[\frac{n-s+1}{2}\right]+1=-s+1$ or -s+2, and $2\left[\frac{n-s+1}{2}\right]-2=n-s-1$ or n-s-2. These give the required

upper estimate for the solution u. For the second part of the theorem, let us notice that when m-2p=s, the corresponding number q is $q=\left[\frac{m-s+1}{2}\right]=p$. Thus, the lower estimate in Theorem 5.4 gives

$$|u_m^j|_{t=\rho}| \ge c\rho^{-m+2q+1} = c\rho^{-s+1}$$

which completes the proof. \square

Acknowledgements. This research is partially supported by the Bulgarian NSC under Grant MM-904/99 and by NSF of Sofia University.

REFERENCES

- 1. Aldashev, S. A. Correctness of multi-dimensional Darboux problems for the wave equation. *Ukr. Math. J.*, **45**, No.9, 1993, 1456-1464.
- 2. Aldashev, S. A. On Darboux problems for a class of multidimensional hyperbolic equations. *Differ. Equations*, **34**, 1998, 65-69.
- 3. Aldashev, S. A. Some problems for a multidimensional hyperbolic integro-differential equation. *Ukrainian Math. J.*. **52**, no. 5, 2000, 590-595.
- Aziz, A. K., M. Schneider. Frankl-Morawetz problems in R³. SIAM J. Math. Anal., 10, 1979, 913-921.
- Bazarbekov, Ar. B., Ak. B. Bazarbekov. Goursat and Darboux problems for the two-dimensional wave equation. I. Differ. Equations, 30, 1994, 741-748.
- Bitsadze, A. V. Some classes of partial differential equations. Gordon and Breach Science Publishers, New York, 1988.
- Copson, E. T. On the Riemann-Green function. J. Rath. Mech. Anal., 1, 1958, 324-348.
- 8. Edmunds, D. E., N. I. Popivanov. A nonlocal regularization of some over-determined boundary value problems. *I. SIAM J. Math. Anal.*, **29**, No1, 1998, 85-105.
- Erdélyi, A., W. Magnus, F. Oberhettinger, F. Tricomi. Higher Transcendental Functions. 1. McGraw-Hill, New York, 1953.
- 10. Garabedian, P. R. Partial differential equations with more than two variables in the complex domain. J. Math. Mech., 9, 1960, 241-271.
- Grammatikopoulos, M. K., T. D. Hristov, N. I. Popivanov. On the singularities of 3-D Protter's problem for the wave equation. *Electron. J. Diff. Eqns.* [online], No.01, 2001, 1-26. Available from: http://ejde.math.swt.cdu/volumes/2001/01/.

- Grammatikopoulos, M. K., T. D. Hristov, N. I. Popivanov. Singular solutions to Protter's problem for the 3-D wave equation involving lower order terms. *Electron.* J. Diff. Eqns. [online], 2003, No. 03, 1-31. Available from: http://ejde.math.swt.edu/volumes/2003/03/.
- Grammatikopoulos, M. K., N. I. Popivanov, T. P. Popov. New singular solutions of Protter's problem for the 3 - D wave equation. Abstract and Applied Analysis, 2003 (in print).
- Grisvard, P. Singularities in BVP. Research in Applied Math., Masson, Paris, Springer-Verlag, Berlin, 1992.
- Hörmander, L. The Analysis of Linear Partial Differential Operators III. Springer-Verlag, Berlin-Heidelberg-New York-Tokyo, 1985.
- Jeon, Jong Duek, Khe Kan Cher, Ji Hyun Park, Yong Hee Jeon, Jong Bae Choi. Protter's conjugate boundary value problems for the two dimensional wave equation. J. Korean. Math. Soc., 33, 1996, 857-863.
- Karatoprakliev, G. D. Uniqueness of solutions of certain boundary-value problems for equations of mixed type and hyperbolic equations in space. *Differ. Equations*, 18, 1982, 49-53.
- 18. Kharibegashvili, S. On the solvability of a spatial problem of Darboux type for the wave equation. Georgian Math. J., 2, 1995, 385-394.
- Khe, Kan Cher. Darboux-Protter problems for the multidimensional wave equation in the class of unbounded functions. Math. Notices of Jacutsk State Univ., 2, 1995, 105-109.
- Khe, Kan Cher. Nontrivial solutions of some homogeneous boundary value problems for a many-dimentional hyperbolic Euler-Poisson-Darboux equation in an unbounded domain. Differ. Equations, 34, 1998, 139-142.
- 21. Kondratiev, V. A. Boundary value problems for elliptic equations in domains with conical and angular points. *Trans. Moscow Math. Soc.*, **16**, 1966, 227-313.
- Marichev, O. I. Handbook of Integral Transforms of Higher Transcendental Functions: Theory and Algorithmic Tables. Ellis Horwood Series in Mathematics and its Applications, Ellis Horwood Limited (Chichester) (a division of John Wiley & Sons), 1983.
- Nakhushev, A. M. Criteria for continuity of the gradient of the solution to the Darboux problem for the Gellerstedt equation. Differ. Equations, 28, 1992, 1445-1457.
- Popivanov, N. I., M. Schneider. The Darboux problems in ℝ³ for a class of degenerated hyperbolic equations. Comptes Rend. de l'Acad. Bulg. Sci., 41, No.11, 1988, 537-578.
- Popivanov, N. I., M. Schneider. On M. H. Protter problems for the wave equation in ℝ³. J. Math. Anal. Appl., 194, No.1, 1995, 50-77.
- Popivanov, N. I., T. P. Popov. Exact Behavior of Singularities of Protter's Problem for the 3-D Wave Equation. In: Herzberger, J.(ed), "Inclusion Methods for Nonlinear Problems" with Applications in Engineering, Economics and Physics, Computing [Suppl.], 16, 2002, 213-236.
- Protter, M. H. New boundary value problems for the wave equation and equations of mixed type. J. Rat. Mech. Anal., 3, 1954, 435-446.
- 28. Samko, S. G., A. A. Kilbas, O.I. Marichev. Fractional Integrals and Derivatives: Theory and Applications. Gordon and Breach, New York, 1993.

 Tong, Kwang-Chang. On a boundary-value problem for the wave equation. Science Record, New Series, 1, No.1, 1957, 1-3.

Received on April 4, 2003

Faculty of Mathematics and Informatics "St. Kl. Ohridski" University of Sofia 5, J. Bourchier blvd., 1164 Sofia BULGARIA E-mail: nedyu@fmi.uni-sofia.bg topopover@hotmail.comv